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ON DYNAMICAL SYSTEM OF CONFLICT WITH FAIR  
REDISTRIBUTION OF VITAL RESOURCES

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We discuss the properties of dynamical system of conflict (DSC). The notion of DSC was introduced in works [6, 7]) for modelling the alternative interaction between opponents.

Let the probability measures  $\mu, \nu$  describe the starting distribution of the vital resource space  $\Omega$  for a pair of opponents  $A, B$ .

The problem is to find the **law** of conflict interaction  $\star$  between  $A, B$  which ensures the compromise redistribution of  $\Omega$ .

Assume that the evolution changes of  $\mu(t), \nu(t)$  are governed by the following nonlinear law of conflict dynamic (cf. with [1, 5]):

$$\dot{\mu} = \frac{\mu\Theta - \tau}{z}, \quad \dot{\nu} = \frac{\nu\Theta - \tau}{z},$$

where  $\Theta = \Theta(\mu, \nu)$  is a positive quadratic form which fixes the so called conflict exponent describing the

**global conflict interaction**

between opponents  $A, B$  and  $\tau = \tau(\mu, \nu)$  has sense of the **occupation**

exponent (the local interaction).

The meanings of  $\tau$  at each moment of time show the values of presence of opponents  $A, B$  on the opposite territory. The denominator  $z$  ensures that measures  $\mu(t), \nu(t)$  are probability for all  $t > 0$ .

We prove that under the appropriate construction of  $\Theta$  and  $\tau$  there exists the  $\omega$ -limit state  $\{\mu^\infty, \nu^\infty\}$  which corresponds to the

**fair redistribution**

of the vital resources space  $\Omega$  between opponents  $A, B$ . The fair redistribution means that the opponents  $A, B$  reach the

**compromise, equilibrium state.**

That is the limit measures  $\mu^\infty, \nu^\infty$  coincide with the normalized components of the classic Jordan decomposition of the signed measure  $\omega = \mu - \nu = \omega_+ - \omega_-$ , i.e.

$$\mu^\infty = \mu_+ := \frac{\omega_+}{\omega_+(\Omega)}, \quad \nu^\infty = \nu_- := \frac{\omega_-}{\omega_-(\Omega)}.$$

In what follows we come to the discrete time  $t = N = 0, 1, \dots$  and use the system of difference equations:

$$\begin{cases} \mu^{N+1}(E) = \mu^N(E) + \mu^N(E)\Theta^N - \tau^N(E), \\ \nu^{N+1}(E) = \nu^N(E) + \nu^N(E)\Theta^N - \tau^N(E), \end{cases} E \in \mathcal{B}, \quad (1)$$

where we omit normalization denominators.

In [11] it was proved (see also [6, 7, 12]) that each trajectory  $\{\mu^N, \nu^N\}$ ,  $N \geq 1$  starting with any couple of probability measures  $\mu, \nu \in \mathcal{M}_1^+(\Omega)$ ,  $\mu \neq \nu$ , converges in the weak sense to a limit fixed point  $\{\mu^\infty, \nu^\infty\}$ . This point creates an equilibrium state for the system and is compromise in the sense  $\mu^\infty \perp \nu^\infty$ . Moreover, for each dynamical system of conflict  $\{\Omega, \mathcal{M}_1^+(\Omega), \star\}$  given by (1), there exists the limit  $\omega$ -set  $\Gamma^\infty$  [13]. It is attractor consisting all couples of mutually singular measures from  $\mathcal{M}_1^+(\Omega)$ . Thus,

$$\Gamma^\infty = \{\{\mu^\infty, \nu^\infty\} \mid \mu^\infty, \nu^\infty \in \mathcal{M}_1^+(\Omega), \mu^\infty \perp \nu^\infty\}.$$

In the simplest situation the dynamical system of conflict can be written in terms of coordinates of stochastic vectors  $p, r \in \mathbf{R}_+^n, n \geq 2$  corresponding to the opponent sides:

$$p_i^{N+1} = 1/z^N(p_i^N \Theta^N - \tau_i^N), \quad r_i^{N+1} = 1/z^N(r_i^N \Theta^N - \tau_i^N), \quad i = 1, \dots, n.$$

Here  $\Theta^N = (p^N, r^N)$  be the inner product between vectors  $p^N, r^N$  and  $\tau_i^N = p_i^N r_i^N$ . We claim that each trajectory  $\{p^N, r^N\}_{N=0}^\infty$  starting with a couple of stochastic vectors  $\{p^0 = p, r^0 = r\}$ ,  $p \neq r$  converges with  $N \rightarrow \infty$  to a fixed point  $\{p^\infty, r^\infty\}$  which creates a compromise state,  $p^\infty \perp r^\infty$ . This state is uniquely determined by the starting couple  $\{p, r\}$  and has an explicit coordinate representation:

$$p_i^\infty = \frac{d_i}{D} > 0, \quad i \in N_+, \quad r_k^\infty = -\frac{d_k}{D} > 0, \quad k \in N_-, \quad D = 1/2 \sum_{i=1}^n |d_i|, \quad (2)$$

$$p_i^\infty = 0, \quad i \notin N_+, \quad r_k^\infty = 0, \quad k \notin N_-,$$

where

$$d_i = p_i - r_i, \quad N_+ = \{i : d_i > 0\}, \quad N_- = \{k : d_k < 0\}.$$

In [10, 9] we generalized above constructions for the cases of piecewise uniformly distributed measures, self-similar, and similar structure measures.

In the most general situation we obtain the similar results in terms of abstract measures (for detail see [11, 12]).

We will deal with DSC  $\{\Omega, \mathcal{M}_{1,ac}^+(\Omega), \star\}$  of natural conflict (here  $\mathcal{M}_{1,ac}^+$  denotes a class of the absolutely continuous measures). A term "natural" means that conflict composition  $\star$  is defined by a fixed law of the conflict interaction between opponents and their strategies do not change during the time evolution.

Let us consider an abstract variant of DSC at discrete time

$$\mu^{N+1} = \mu^N \star \nu^N, \quad \nu^{N+1} = \nu^N \star \mu^N, \quad N = 0, 1, \dots$$

Their state trajectories

$$\left\{ \begin{array}{c} \mu^N \\ \nu^N \end{array} \right\} \xrightarrow{\star} \left\{ \begin{array}{c} \mu^{N+1} \\ \nu^{N+1} \end{array} \right\}, \quad N = 0, 1, \dots \quad (3)$$

are governed by the following *law of conflict dynamic*:

$$\left\{ \begin{array}{l} \mu^{N+1}(E) = \frac{1}{z^N}[\mu^N(E)(\Theta^N + 1) - \tau^N(E)], \\ \nu^{N+1}(E) = \frac{1}{z^N}[\nu^N(E)(\Theta^N + 1) - \tau^N(E)], \end{array} \right. \quad E \in \mathcal{B}, \quad (4)$$

where measures  $\mu^0 = \mu$ ,  $\nu^0 = \nu$  correspond to an initial state. The conflict exponent  $\Theta^N$  in (4) is defined as

$$\Theta^N = \int_{\Omega} \int_{\Omega} \mathcal{K}(x, y) \varphi^N(x) \psi^N(y) dx dy,$$

where  $\mathcal{K}(x, y)$  denotes a kernel of some positive bounded operator  $K$  in  $L^2(\Omega, d\lambda)$  and

$$\varphi^N(x) = \sqrt{\rho^N(x)}, \quad \psi^N(x) = \sqrt{\sigma^N(x)},$$

where  $\rho^N(x), \sigma^N(x)$  are the Radon-Nikodym derivatives of  $\mu^N, \nu^N$  with respect to  $\lambda$ . Thus,

$$\Theta_N = (K\varphi^N, \psi^N)_{L^2(\Omega, d\lambda)},$$

Further,  $\tau^N$  in (4) stands for the occupation measure. Its values characterize the presence of opponents on opposite territories. By definition,

$$\tau^N(E) = \nu^N(E_+) + \mu^N(E_-), \quad E_+ = E \cap \Omega_+, \quad E_- = E \cap \Omega_-, \quad (5)$$

where  $\Omega = \Omega_- \cup \Omega_+$  corresponds to the Hahn-Jordan decomposition (see [3, 14]) of the starting signed measure  $\omega = \mu - \nu$ . Finally, the normalizing denominator in (4) is defined as

$$z^N = \Theta^N + 1 - W^N, \quad W^N = \mu^N(\Omega_-) + \nu^N(\Omega_+).$$

It is easy to see that all measures  $\mu^N, \nu^N$ ,  $N \geq 1$  in (4) are absolutely continuous and probability, i.e.,  $\mu^N, \nu^N \in \mathcal{M}_{1,\text{ac}}^+(\Omega)$ .

The DSC defined by (4) has two separate sets of fixed points. The first set contains all couples of identical measures from  $\mathcal{M}_{1,\text{ac}}^+(\Omega)$ . Indeed, if  $\mu = \nu$ , then  $\Theta^N = \text{const}$  for all  $N$  and  $\mu^N(E) = \mu(E) = \nu^N(E) = \nu(E)$  for each  $E \in \mathcal{B}$ . The second set is composed from measures  $\mu, \nu \in \mathcal{M}_{1,\text{ac}}^+(\Omega)$  which are orthogonal,  $\mu \perp \nu$ . In this case  $\tau^N = 0 = W^N$  and  $\Theta^N + 1 = z^N$  for all  $N$ . Due to (4) we find that  $\mu^N = \mu$ ,  $\nu^N = \nu$ .

In all other cases, when the starting measures are different,  $\mu \neq \nu$ , and mutually non-singular the following theorem is true.

**Theorem** Let  $\{\Omega, \mathcal{M}_{1,\text{ac}}^+(\Omega), \star\}$  be a DSC generated by the system of difference equations (4). Then each its trajectory (3) starting with a couple of probability measures  $\mu^0 = \mu, \nu^0 = \nu \in \mathcal{M}_{1,\text{ac}}^+(\Omega)$ ,  $\mu \neq \nu$  converges to the fixed point corresponding to the limit state  $\{\mu^\infty, \nu^\infty\}$  with

$$\mu^\infty(E) = \lim_{N \rightarrow \infty} \mu^N(E), \quad \nu^\infty(E) = \lim_{N \rightarrow \infty} \nu^N(E), \quad E \in \mathcal{B}, \quad (6)$$

where

$$\begin{aligned} \mu^\infty(E) &= \frac{\mu(E_+) - \nu(E_+)}{D} = \mu_+(E), \\ \nu^\infty(E) &= -\frac{\mu(E_-) - \nu(E_-)}{D} = \nu_-(E) \end{aligned} \quad (7)$$

with  $\mu_+, \nu_-$  defined as the Hahn normalized components of  $\omega = \mu - \nu$ .

In (7)  $D = 1/2 \int_\Omega |\rho(x) - \sigma(x)| dx$  stands for the total difference between measures  $\mu, \nu$ .

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